

Hand in until Tuesday, July 12, 09:45

## Exercise sheet 5

## Exercise 1 (5 Points)

If an operator A is an eigenvector of the Liouville operator, i.e. LA = aA, show that the projection operator

$$P = |A\rangle \frac{1}{\langle A|A\rangle} \langle A| \tag{1}$$

commutes with the Liouville operator.

## Exercise 2 (40 Points)

Consider the relaxation of non–interacting spins in a magnetic field due to lattice modulations. This problem can be formulated as spins interacting with a magnetic field and a bosonic bath which represents the phononic excitations. The Hamiltonian that describes this problem is given by

$$H = H_s + H_b + H_{sb}$$
  
=  $-g\mu_B BS^z + \sum_{\mathbf{k}} \omega_{\mathbf{k}} b_{\mathbf{k}}^{\dagger} b_{\mathbf{k}} + \frac{1}{2} \sum_{\mathbf{k}} \left( \lambda_{\mathbf{k}} b_{\mathbf{k}} + \lambda_{\mathbf{k}}^* b_{\mathbf{k}}^{\dagger} \right) S^x .$  (2)

In the first term, g is the Landé factor,  $\mu_B$  the Bohr magneton, *B* the magnetic field, and  $S^z$  the *z* component of the spin–1/2 operator. In the second term,  $\omega_k$  is the dispersion of the phonons, and  $b^{\dagger}$ , *b* are bosonic creation, annihilation operators obeying the commutation relations

$$[b_{{f k}},b_{{f k}'}]=0=[b_{{f k}}^{\dagger},b_{{f k}'}^{\dagger}] \ \ \ ext{and} \ \ \ [b_{{f k}},b_{{f k}'}^{\dagger}]=\delta_{{f k},{f k}'}\,.$$

Lastly, in the third term,  $\lambda_k$  is the coupling strength of the bosonic bath to the spins, and  $S^x$  is the *x* component of the spin operator which can be written in terms of raising(*S*<sup>+</sup>), lowering(*S*<sup>-</sup>) operators as  $S^x = \frac{1}{2}(S^+ + S^-)$ . In the rotating wave approximation, the interaction term between spin and phonons reduces to

$$H_{sb} \approx \sum_{\mathbf{k}} \left( \lambda_{\mathbf{k}} b_{\mathbf{k}} S^{-} + \lambda_{\mathbf{k}}^{*} b_{\mathbf{k}}^{\dagger} S^{+} \right) .$$
(3)

Furthermore, for weak  $\lambda_k$ ,  $H_{sb}$  can be considered as a perturbation to the non interacting system, viz.

$$H = H_0 + H_1$$
, with  $H_0 = H_s + H_b$ , and  $H_1 = H_{sb}$ .

- Show that in the unperturbed system (λ<sub>k</sub> = 0), the spin operators S<sup>z</sup>, S<sup>±</sup> as well as, the bosonic b<sub>k</sub>, b<sup>†</sup><sub>k</sub> are eigenvectors of the Liouville operator L<sub>0</sub>.
- In the unperturbed system, calculate the static susceptibility χ<sub>0</sub><sup>zz</sup> for the S<sup>z</sup> operator, and the imaginary part of the dynamical susceptibility ℑ[χ<sub>0</sub><sup>bb</sup>(ω)] with

$$\Im[\chi_0^{bb}(\omega)] = \int_{-\infty}^{\infty} dt e^{i\omega t} \langle [b_{\mathbf{k}}(t), b_{\mathbf{k}'}^{\dagger}] \rangle_0 .$$
<sup>(4)</sup>

- Show that the operator  $L_1S^z$  is orthogonal to  $S^z$  with respect to the unperturbed system.
- Show that in second order of the coupling  $\lambda$ , the memory function for the  $S^z$  operator is given by

$$\mathcal{M}^{zz}(z) \simeq \langle QL_1 S^z | \frac{1}{z - L_0} QL_1 S^z \rangle_0 + \mathcal{O}(\lambda^3) .$$
(5)

• Calculate the memory function  $M^{zz}(z)$  to second order in the interaction coupling  $\lambda$ .

## Exercise 3 (10 Points)

In the long time approximation, an observable A is considered to be slowly varying in some characteristic time scale  $\tau$ . In that case, the memory function  $M_{AA}(z)$  can be considered as frequency independent at the origin of the complex plane. Describe the relaxation function  $\Phi_{AA}(t)$  for the observable A.