

Hand in until Tuesday, June 21, 09:45

Exercise sheet 3

🛱 Exercise 1 (15 Points)

In the one–dimensional tight binding model, spinless fermions hop from one site of a ring to the other with a hopping amplitude t, see Fig. 1. Assuming also a finite chemical potential μ , the system is described by the Hamiltonian

$$H = -t \sum_{l=0}^{N-1} \left(c_l^{\dagger} c_{l+1} + c_{l+1}^{\dagger} c_l \right) - \mu \sum_{l=0}^{N-1} c_l^{\dagger} c_l .$$
⁽¹⁾

Here, N is the number of lattice sites, and $c_l^{\dagger}(c_l)$ are creation(annihilation) operators obeying the usual fermionic anticommutation relations

$$\{c_l, c_{l'}^{\dagger}\} = \delta_{ll'}, \quad \text{and} \quad \{c_l, c_{l'}\} = 0 = \{c_l^{\dagger}, c_{l'}^{\dagger}\}.$$
(2)

- By going to the reciprocal space diagonalize Hamiltonian (1).
- Calculate the ground state of the system as a function of the chemical potential μ and the bandwidth t.
- Calculate the ground state for *N*/2 fermions.
- Calculate the density–density dynamical susceptibility $\chi_{nn}(q, z)$, defined as*

$$\chi_{nn}(q,z) = -\langle [n(-q), \frac{1}{z-L}n(q)] \rangle , \qquad (3)$$

with *q* the wavevector, $z = \omega + i0^+$ the complex frequency, *L* the Liouville operator, and *n* the density operator

$$n = \frac{1}{N} \sum_{l} c_l^{\dagger} c_l . \tag{4}$$

• Give a graphical representation of the imaginary part of the susceptibility $\chi''_{nn}(q, \omega)$, at zero temperature, as a function of ω , and q, for N/2 particles. What do you conclude? What happens for N particles?

Exercise 2 (40 Points)

Consider a one–dimensional tight binding model on a ring with 2 atoms per unit cell. This results to a model where spinless fermions hop from one site to the other with two alternating hopping amplitudes t_1 , and t_2 , see Fig. 1. The Hamiltonian that describes this system is given by

$$H = -\mu \sum_{l=0}^{N-1} a_{2l}^{\dagger} a_{2l} + b_{2l+1}^{\dagger} b_{2l+1} - t_1 \sum_{l=0}^{N-1} \left(a_{2l}^{\dagger} b_{2l+1} + b_{2l+1}^{\dagger} a_{2l} \right) - t_2 \sum_{l=0}^{N-1} \left(a_{2(l+1)}^{\dagger} b_{2l+1} + b_{2l+1}^{\dagger} a_{2(l+1)} \right) .$$
(5)

^{*} Hint: Think of the action of the Liouville operator on the density operator in the Hamiltonian basis. What about powers of the Liouvillian operator?



Figure 1: (left) A tight binding model on a ring with one atom per unit cell. Fermions hop from site to site with an amplitude t. (right) A tight binding model on a ring with two atoms per unit cell. Fermions hop from site to site with two different hopping amplitudes t_1 and t_2 .

Here, $N = N_s/2$ is the number of unit cells, N_s is the number of sites, μ is the chemical potential, and a_l^{\dagger} , b_l^{\dagger} (a_l , b_l) are creation(annihilation) operators. a and b operators always anticommute with each other while for each species we have the anticommutation relations

$$\{a_l, a_{l'}^{\dagger}\} = \delta_{ll'}, \text{ and } \{a_l, a_{l'}\} = 0 = \{a_l^{\dagger}, a_{l'}^{\dagger}\},$$
(6)

$$\{b_l, b_{l'}^{\dagger}\} = \delta_{ll'}, \text{ and } \{b_l, b_{l'}\} = 0 = \{b_l^{\dagger}, b_{l'}^{\dagger}\}.$$
(7)

Show that Hamiltonian (5) can be written in a matrix form in reciprocal space as

$$H = \sum_{k} \Psi_{k}^{\dagger} \mathcal{H}_{k} \Psi_{k}$$
(8)

where $\Psi_k^{\dagger} = (a_k^{\dagger}, b_k^{\dagger})$, \mathcal{H}_k some 2 × 2 matrix and k is the quantized momentum index. Note that the sum in Eq. (8) runs over the first Brillouin zone.

- Diagonalize Hamiltonian (8) and find the energy dispersion relations.
- Calculate the density-density dynamical susceptibilities $\chi_{aa}(q, z)$, $\chi_{bb}(q, z)$ defined as (see footnote)

$$\chi_{aa}(q,z) = -\langle [n_a(-q), \frac{1}{z-L}n_a(q)] \rangle , \qquad \chi_{bb}(q,z) = -\langle [n_b(-q), \frac{1}{z-L}n_b(q)] \rangle , \tag{9}$$

with q the wavevector, $z = \omega + i0^+$ the complex frequency, L the Liouville operator, and n_a, n_b the density operators

$$n_a = \frac{1}{N} \sum_{l=0}^{N-1} a_{2l}^{\dagger} a_{2l} , \qquad n_b = \frac{1}{N} \sum_{l=0}^{N-1} b_{2l+1}^{\dagger} b_{2l+1} .$$
 (10)

• Give a graphical representation of the imaginary part of the susceptibilities $\chi_{aa}^{\prime\prime}(q,\omega)$, and $\chi_{bb}^{\prime\prime}(q,\omega)$ for $t_1 = -t_2$, at zero temperature, as a function of ω , and q, for N particles. What do you conclude?