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Partitioned Methods for Multifield Problems

Rang, 21.6.2016



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Generalised- α method

Consider

$$\dot{\mathbf{u}} = \mathbf{f}(t, \mathbf{u}), \quad \mathbf{u}(0) = \mathbf{u}_0.$$

generalised- α method:

$$\begin{aligned}\dot{\mathbf{u}}_{n+\alpha_m} &= \mathbf{f}(t_{n+\alpha_f}, \mathbf{u}_{n+\alpha_f}), \\ \mathbf{u}_{n+1} &= \mathbf{u}_n + \tau \dot{\mathbf{u}}_n + \tau \gamma (\dot{\mathbf{u}}_{n+1} - \dot{\mathbf{u}}_n), \\ \dot{\mathbf{u}}_{n+\alpha_m} &= \dot{\mathbf{u}}_n + \alpha_m (\dot{\mathbf{u}}_{n+1} - \dot{\mathbf{u}}_n), \\ \mathbf{u}_{n+\alpha_f} &= \mathbf{u}_n + \alpha_f (\mathbf{u}_{n+1} - \mathbf{u}_n).\end{aligned}$$

Formulation as onestep method

$$\mathbf{u}_{n+1} = \mathbf{u}_n + \tau \left(1 - \frac{\gamma}{\alpha_m} \right) \dot{\mathbf{u}}_n + \frac{\tau\gamma}{\alpha_m} \mathbf{f}_{n+\alpha_f},$$
$$\dot{\mathbf{u}}_{n+1} = \frac{1}{\tau\gamma} (\mathbf{u}_{n+1} - \mathbf{u}_n - \tau(1 - \gamma)\dot{\mathbf{u}}_n),$$

if $\alpha_m \neq 0$.

- $\dot{\mathbf{u}}_0$ can be computed from the ODE.
- Consistency order 2, if $\alpha_m = \alpha_f$ and $\gamma = 1/2$.
- zero-stability and convergency, if $\alpha_m > 1/2$.

Formulation as multistep method

Formulation:

$$\mathbf{u}_{n+1} = \frac{2\alpha_m - 1}{\alpha_m} \mathbf{u}_n - \frac{\alpha_m - 1}{\alpha_m} \mathbf{u}_{n-1} + \frac{\tau(1 - \gamma)}{\alpha_m} \mathbf{f}_{n-1+\alpha_f} + \frac{\tau\gamma}{\alpha_m} \mathbf{f}_{n+\alpha_f}.$$

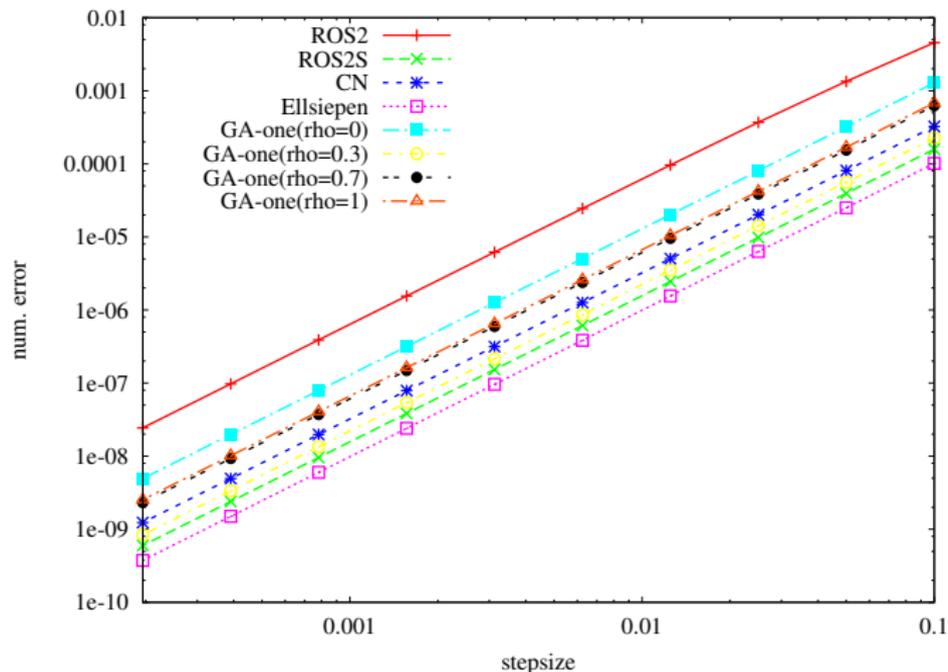
- second order consistency:

$$\gamma = \frac{1}{2} - \alpha_f + \alpha_m, \quad (1)$$

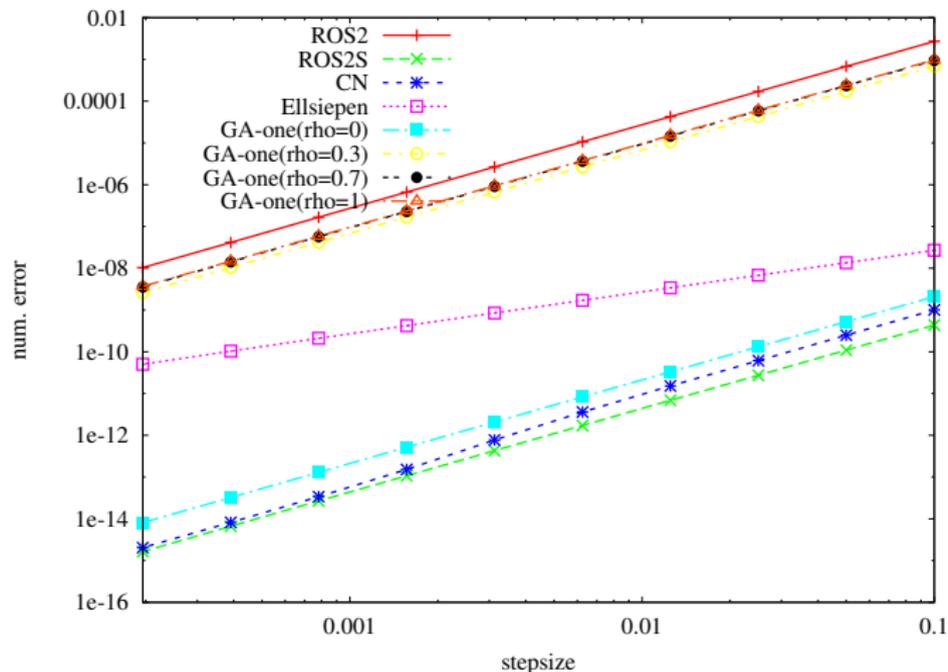
- generalised- α method is convergent of order 2 if $\alpha_m > 1/2$ and (1) hold.
- setting due to stability reasons:

$$\alpha_f = \gamma = \frac{1}{1 + \rho_\infty}, \quad \alpha_m = \frac{3 - \rho_\infty}{2(1 + \rho_\infty)}.$$

Example of Prothero and Robinson ($\lambda = -1$)



Example of Prothero and Robinson ($\lambda = -1.0E+6$)



Comparison

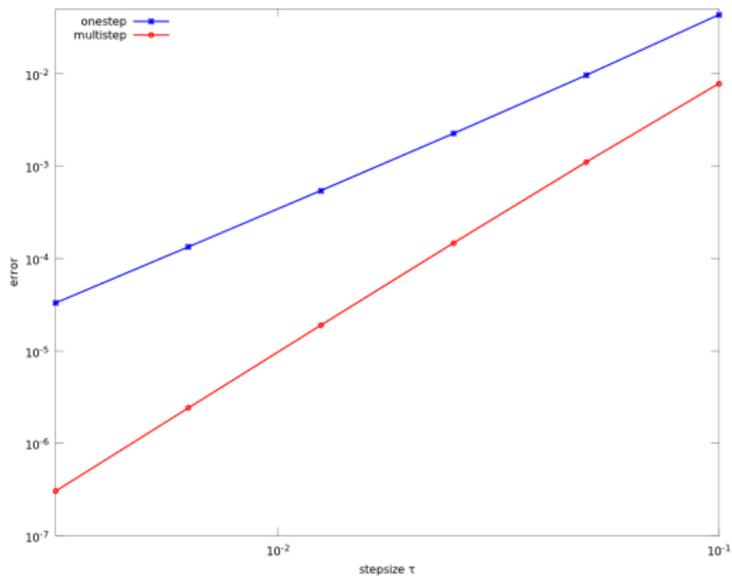
- Let $\alpha_f = 1$, $\alpha_m = 5/6$, $\gamma = 1/3 \rightarrow$ **3rd order for multistep version**
- Let $\tau = 1/(10 \cdot 2^k)$ with $k = 0, \dots, 5$
- Consider the second order ODE

$$\ddot{y}_i = -\frac{y_i}{(y_1^2 + y_2^2)^{3/2}}, \quad i = 1, 2.$$

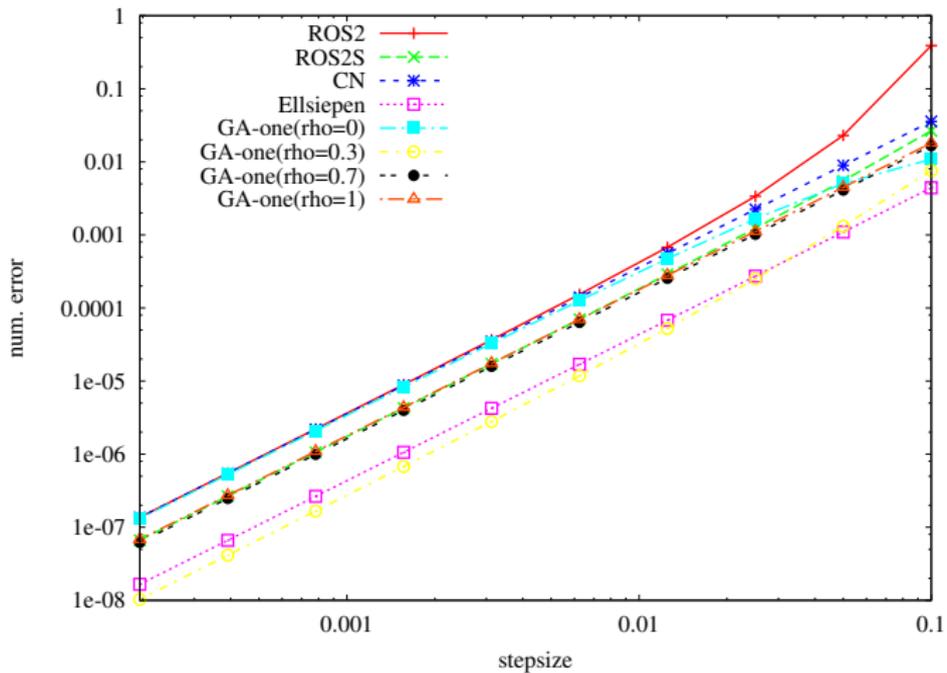
The initial conditions are given by $\mathbf{u}_0 = \left(0, \sqrt{\frac{1+e}{1-e}}, 1-e, 0\right)^T$,
where $e = 1/2$.

- time interval: $[0, 10]$

Kepler's problem



Comparison with other 2nd order methods



Formulation as onestep method

Consider the second order ODE

$$\ddot{\mathbf{u}} = \mathbf{f}(t, \mathbf{u}, \dot{\mathbf{u}}), \quad \mathbf{u}(0) = \mathbf{u}_0, \dot{\mathbf{u}}(0) = \dot{\mathbf{u}}_0. \quad (2)$$

generalised- α method:

$$\mathbf{u}_{n+\alpha_f} = \alpha_f \mathbf{u}_{n+1} + (1 - \alpha_f) \mathbf{u}_n,$$

$$\dot{\mathbf{u}}_{n+\alpha_f} = \alpha_f \dot{\mathbf{u}}_{n+1} + (1 - \alpha_f) \dot{\mathbf{u}}_n,$$

$$\ddot{\mathbf{u}}_{n+\alpha_m} = \alpha_m \ddot{\mathbf{u}}_{n+1} + (1 - \alpha_m) \ddot{\mathbf{u}}_n,$$

$$\mathbf{u}_{n+1} = \mathbf{u}_n + \tau \dot{\mathbf{u}}_n + \tau^2 \left[\left(\frac{1}{2} - \beta \right) \ddot{\mathbf{u}}_n + \beta \ddot{\mathbf{u}}_{n+1} \right],$$

$$\dot{\mathbf{u}}_{n+1} = \dot{\mathbf{u}}_n + \tau [(1 - \gamma) \ddot{\mathbf{u}}_n + \gamma \ddot{\mathbf{u}}_{n+1}],$$

$$\ddot{\mathbf{u}}_{n+\alpha_m} = \mathbf{f}(t_{n+\alpha_f}, \alpha_f \mathbf{u}_{n+1} + (1 - \alpha_f) \mathbf{u}_n, \alpha_f \dot{\mathbf{u}}_{n+1} + (1 - \alpha_f) \dot{\mathbf{u}}_n),$$

where $t_{n+\alpha_f} = t_n + \tau \alpha_f$.

Formulation as onestep method

Abbreviation:

$$\mathbf{f}_{n+\alpha_f} := \mathbf{f}(t_{n+\alpha_f}, \alpha_f \mathbf{u}_{n+1} + (1 - \alpha_f) \mathbf{u}_n, \alpha_f \dot{\mathbf{u}}_{n+1} + (1 - \alpha_f) \dot{\mathbf{u}}_n).$$

Then

$$\mathbf{u}_{n+1} = \mathbf{u}_n + \tau \dot{\mathbf{u}}_n + \tau^2 \left[\left(\frac{1}{2} - \frac{\beta}{\alpha_m} \right) \ddot{\mathbf{u}}_n + \frac{\beta}{\alpha_m} \mathbf{f}_{n+\alpha_f} \right],$$

$$\dot{\mathbf{u}}_{n+1} = \dot{\mathbf{u}}_n + \tau \left[\left(1 - \frac{\gamma}{\alpha_m} \right) \ddot{\mathbf{u}}_n + \frac{\gamma}{\alpha_m} \mathbf{f}_{n+\alpha_f} \right].$$

$$\ddot{\mathbf{u}}_{n+1} = \frac{1}{\alpha_m} [\ddot{\mathbf{u}}_{n+\alpha_m} - (1 - \alpha_m) \ddot{\mathbf{u}}_n] = \frac{1}{\alpha_m} [\mathbf{f}_{n+\alpha_f} - (1 - \alpha_m) \ddot{\mathbf{u}}_n].$$

Formulation as multistep method

Consider

$$M\ddot{\mathbf{u}} + C\dot{\mathbf{u}} + \mathbf{S}(\mathbf{u}) = \mathbf{F}(t), \mathbf{u}(0) = \mathbf{u}_0, \quad \dot{\mathbf{u}}_0 = \mathbf{v}_0.$$

Then

$$M\ddot{\mathbf{u}}_{n+\alpha_m} = \mathbf{F}(t_{n+\alpha_f}) - \mathbf{S}(\alpha_f \mathbf{u}_{n+1} + (1 - \alpha_f) \mathbf{u}_n) - C(\alpha_f \dot{\mathbf{u}}_{n+1} + (1 - \alpha_f) \dot{\mathbf{u}}_n).$$

The formulation as multistep method

Method:

$$\sum_{j=0}^3 [M\alpha_j + \tau C\gamma_j] \mathbf{u}_{n+j} + \tau^2 \sum_{j=0}^2 \delta_j [\mathbf{S}_{n+j+\alpha_f} - \mathbf{F}(t_{n+j+\alpha_f})] = 0,$$

where

$$\alpha_0 = 1 - \alpha_m, \quad \alpha_1 = 3\alpha_m - 2, \quad \alpha_2 = 1 - 3\alpha_m, \quad \alpha_3 = \alpha_m,$$

$$\gamma_0 = (1 - \alpha_f)(\gamma - 1), \quad \gamma_1 = 1 - 2\alpha_f - 2\gamma + 3\gamma\alpha_f, \quad \gamma_2 = \alpha_f + \gamma - 3\gamma\alpha_f,$$

$$\delta_0 = \frac{1}{2} + \beta - \gamma, \quad \delta_1 = \frac{1}{2} - 2\beta + \gamma, \quad \delta_2 = \beta$$

and

$$\mathbf{F}_{n+j-\alpha_f} = \mathbf{F}(\alpha_f t_{n+j+1} + (1 - \alpha_f)t_{n+j}) = \mathbf{F}(t_{n+j} + \alpha_f \tau)$$

$$\mathbf{S}_{n+j+\alpha_f} = \alpha_f \mathbf{S}(\mathbf{u}_{n+j+1}) + (1 - \alpha_f) \mathbf{S}(\mathbf{u}_{n+j}).$$

Properties

- consistency order 2: $\gamma = \frac{1}{2} + \alpha_m - \alpha_f$.
- zero-stable and convergence: $\alpha_m \geq 1/2$, $\alpha_f \leq 1/2$ and $\gamma \leq 1/2$.
- for stability reasons:

$$\beta = \frac{(1 + \alpha_m - \alpha_f)^2}{4}, \alpha_f = \frac{1}{1 + \rho_\infty}, \alpha_m = \frac{2 - \rho_\infty}{1 + \rho_\infty}$$

Comparison

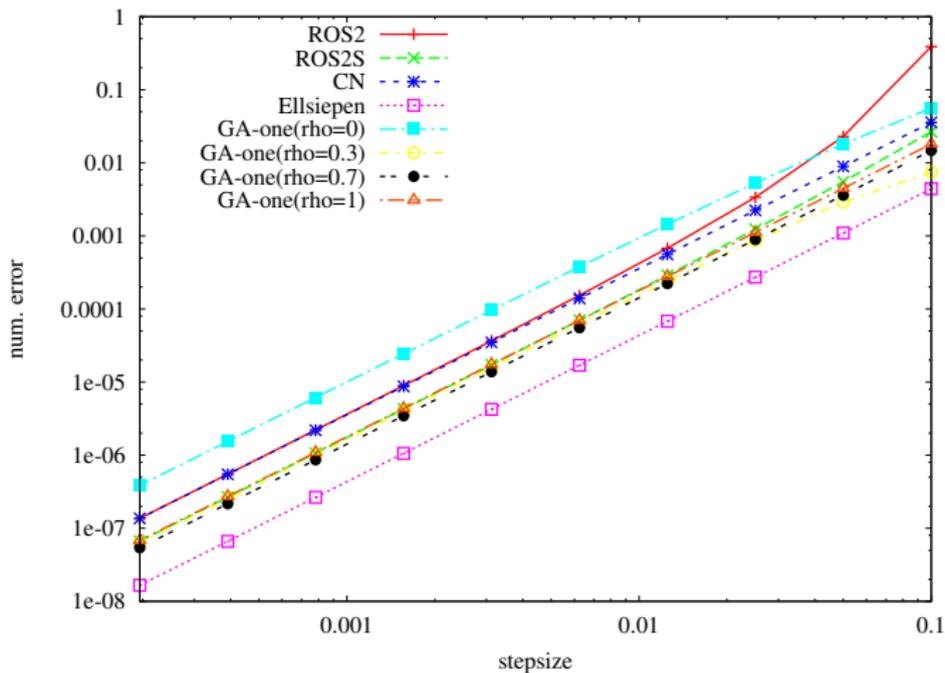
- Let $\tau = 1/(10 \cdot 2^k)$ with $k = 0, \dots, 5$
- Consider the second order ODE

$$\ddot{y}_i = -\frac{y_i}{(y_1^2 + y_2^2)^{3/2}}, \quad i = 1, 2.$$

The initial conditions are given by $\mathbf{u}_0 = \left(0, \sqrt{\frac{1+e}{1-e}}, 1-e, 0\right)^T$,
where $e = 1/2$.

- time interval: $[0, 10]$

Comparison with other 2nd order methods



Adaptivity

- Idea: use PI-controller from Gustafsson et. al. (1988)
- suggest a new timestep size with solutions of order p and $p - 1$.
- use approximation of generalised- α method as a second order approximation.
- use as the second solution with order 1 backward Euler method.
- next timestep size τ_{n+1} :

$$\tau_{n+1} = \rho_{adapt} \frac{\tau_n^2}{\tau_{n-1}} \left(\frac{TOL \cdot r_{n-1}}{r_n^2} \right)^{1/2},$$

where $\rho_{adapt} \in (0, 1]$ is a safety factor, $TOL > 0$ is a given tolerance, and $r_{n+1} := \|\mathbf{u}_{n+1} - \hat{\mathbf{u}}_{n+1}\|$.

Adaptivity for onestep versions

- Compute $(\mathbf{u}_{n+1}, \dot{\mathbf{u}}_{n+1})^\top$ with generalised- α method.
- Compute second solution with the backward Euler method and use $\dot{\mathbf{u}}_{n+1}$ as approximation for $\mathbf{f}(t_{n+1}, \mathbf{u}_{n+1})$, i. e. $\hat{\mathbf{u}}_{n+1} = \mathbf{u}_n + \tau_n \dot{\mathbf{u}}_{n+1}$.
- Compute the numerical error with r_{n+1} and approximate the new timestep length τ_{n+1} .
- If the numerical error is smaller than the given tolerance the timestep is accepted, otherwise it is rejected and has to be recomputed with the new timestep length τ_{n+1} .

The chemical reaction E5

Let $A = 7,89 \times 10^{-10}$, $B = 1,1 \times 10^7$, $C = 1,13 \times 10^3$, and $M = 10^6$.
Consider

$$\dot{u}_1 = -Au_1 - Bu_1u_3,$$

$$\dot{u}_2 = Au_1 - MCu_2u_3,$$

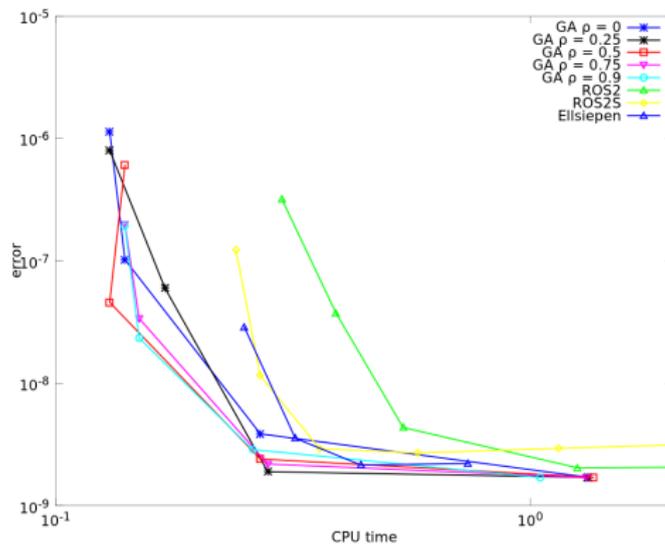
$$\dot{u}_3 = Au_1 - Bu_1u_3 - MCu_2u_3 + Cu_4,$$

$$\dot{u}_4 = Bu_1u_3 - Cu_4$$

with the initial conditions $u_1(0) = 1,76 \times 10^{-3}$ and $u_i(0) = 0$,
 $i \in \{2, 3, 4\}$.

The equations should be solved in the time interval $[0, 10^{13}]$.

The chemical reaction E5



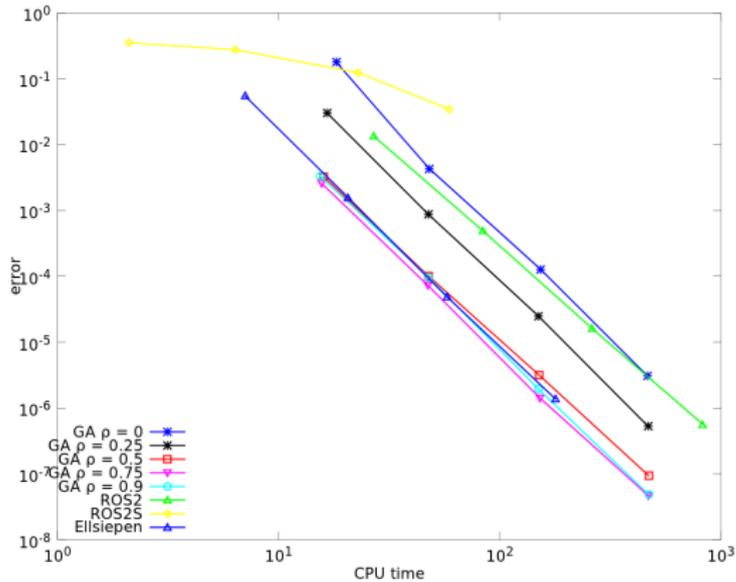
Kepler's problem

Consider the second order ODE

$$\ddot{y}_i = -\frac{y_i}{(y_1^2 + y_2^2)^{3/2}}, \quad i = 1, 2.$$

The initial conditions are given by $\mathbf{u}_0 = \left(0, \sqrt{\frac{1+e}{1-e}}, 1-e, 0\right)^T$, where $e \in [0, 1)$ is a given parameter. In our numerical example we choose $e = 1/2$. We solve the problem in the interval $[0, 20000]$

Kepler's problem



Model problem

Damped spring mass system:

$$(1 - \alpha)\ddot{u}^s + \omega^2 u^s = -\alpha\ddot{u}^f - 2\xi\omega\dot{u}^f$$

with

$$\alpha = \frac{m^f}{m}, \quad 1 - \alpha = \frac{m^s}{m}, \quad m^f + m^s = m,$$
$$\omega = \sqrt{k/m}, \quad \xi = \frac{c}{2\sqrt{km}},$$

Initial conditions:

$$u(0) = u_0, \quad \dot{u}(0) = \dot{u}_0.$$

Model problem

- m ... total mass of the system
- m^f ... mass of the domain f
- m^s ... mass of the domain s
- k ... spring stiffness
- c ... a damping coefficient
- ω ... natural frequency of the system
- ξ ... a dimensionless damping ration

Analytical solution

Consider the case $u := u^f = u^s$. Then

$$\ddot{u} + 2\xi\omega\dot{u} + \omega^2u = 0.$$

- u ... displacement of the total mass
- \dot{u} ... velocity
- \ddot{u} ... acceleration.

Analytical solution

- Ansatz

$$u(t) = c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t},$$

where $\lambda_1, \lambda_2 \in \mathbb{C}$.

- System is undercritically damped, i. e. $\xi < 1$.
- λ_1 and λ_2 can be determined by solving the quadratic problem

$$\lambda^2 + 2\xi\omega\lambda + \omega^2 = 0,$$

- solution:

$$\lambda_{1,2} = -\xi\omega \pm i\omega_D,$$

where $\omega_D = \omega\sqrt{1 - \xi^2}$.

Analytical solution

Then we have

$$u(t) = \exp(-\xi\omega t) [c_1 \exp(-i\omega_D t) + c_2 \exp(i\omega_D t)].$$

First initial condition:

$$u_0 = u(0) = c_1 + c_2.$$

Second initial condition:

$$\dot{u}_0 = \dot{u}(0) = -\xi\omega u_0 + i\omega_D(c_2 - c_1).$$

Analytical solution

$$u(t) = \exp(-\xi\omega t) \left[u_0 \cos(\omega_D t) + \frac{\dot{u}_0 + \xi\omega u_0}{\omega} \sin(\omega_D t) \right].$$

Numerical solution

Idea:

- Discretise LHS with the generalised- α method for second order problems
- Discretise RHS with the generalised- α method for first order problems.

Onestep formulation:

$$(1 - \alpha)\ddot{u}_{n+\alpha_m^s}^s + \omega^2 u_{n+\alpha_f^s}^s = -\alpha\ddot{u}_{n+\alpha_m^f}^f - \xi\omega\dot{u}_{n+\alpha_f^f}^f$$

$$\dot{u}_{n+1}^s = \dot{u}_{n+1}^f =: \dot{u}_{n+1}^I$$

$$\dot{u}_n^s = \dot{u}_n^f =: \dot{u}_n^I.$$

\dot{u}_{n+1}^I and \dot{u}_n^I represent the interface velocities at time t_{n+1} and t_n

Numerical solution

Insert formulas for $\ddot{u}_{n+\alpha_m^s}^s$, $\ddot{u}_{n+\alpha_m^f}^f$, and $\dot{u}_{n+\alpha_f^f}^f$:

$$\begin{aligned} & (1 - \alpha) [\alpha_m^s \ddot{u}_{n+1}^s + (1 - \alpha_m^s) \ddot{u}_n^s] + \omega^2 [\alpha_f^s u_{n+1}^s + (1 - \alpha_f^s) u_n^s] \\ & = -\alpha [\alpha_m^f \ddot{u}_{n+1}^f + (1 - \alpha_m^f) \ddot{u}_n^f] - 2\xi\omega [\alpha_f^f \dot{u}_{n+1}^f + (1 - \alpha_f^f) \dot{u}_n^f]. \end{aligned}$$

Moreover we have

$$\begin{aligned} u_{n+1}^s &= u_n^s + \tau \dot{u}_n^s + \tau^2 \left[\left(\frac{1}{2} - \beta^s \right) \ddot{u}_n^s + \beta_s \ddot{u}_{n+1}^s \right], \\ \dot{u}_{n+1}^f &= \dot{u}_n^f + \tau [\gamma^f \ddot{u}_{n+1}^f + (1 - \gamma^f) \ddot{u}_n^f], \\ \dot{u}_{n+1}^s &= \dot{u}_n^s + \tau [\gamma^s \ddot{u}_{n+1}^s + (1 - \gamma^s) \ddot{u}_n^s]. \end{aligned}$$

Numerical solution

Then $A_1 \mathbf{v}_{n+1} = A_2 \mathbf{v}_n$, where

$$A_1 = \begin{pmatrix} \omega^2 \alpha_f^s & 2 \frac{\xi \omega}{\tau} \alpha_f^f & \frac{(1-\alpha) \alpha_m^s}{\tau^2} & \frac{\alpha \alpha_m^f}{\tau^2} \\ 1 & 0 & -\beta^s & 0 \\ 0 & \frac{1}{\tau} & 0 & -\frac{\gamma^f}{\tau} \\ 0 & \frac{1}{\tau} & -\frac{\gamma^s}{\tau} & 0 \end{pmatrix},$$

$$A_2 = \begin{pmatrix} -\omega^2 (1 - \alpha_f^s) & -2 \frac{\xi \omega}{\tau} (1 - \alpha_f^f) & \frac{(1-\alpha)(1-\alpha_m^s)}{\tau^2} & -\frac{\alpha(1-\alpha_m^f)}{\tau^2} \\ 1 & 1 & \frac{1}{2} - \beta^s & 0 \\ 0 & \frac{1}{\tau} & 0 & \frac{1-\gamma^f}{\tau} \\ 0 & \frac{1}{\tau} & \frac{1-\gamma^s}{\tau} & 0 \end{pmatrix}$$

$$\mathbf{v}_n = (u_n^s, \tau \dot{u}_n^f, \tau^2 \ddot{u}_n^s, \tau^2 \ddot{u}_n^f)^\top.$$

Convergence order

- Write our problem as $\mathbf{v}_{n+1} = A_1^{-1} A_2 \mathbf{v}_n$.
- Expand this expression in a Taylor series

$$u_{n+1}^s = u_n^s + \tau \dot{u}_n^s - \frac{\tau^2}{2[\alpha \alpha_m^f \gamma^s + (1 - \alpha) \alpha_m^s \gamma^f]} [2\beta^s \alpha \alpha_m^f (\ddot{u}_s - \ddot{u}_f) - \alpha \alpha_m^f \gamma^s \ddot{u}_s].$$

Numerical solution with interpolation

- **Problem:** LHS and RHS are taken at different times.
- Let

$$F_{n+\alpha_f^s}^s = (1 - \alpha) \ddot{u}_{n+\alpha_m^s}^s + \omega^2 u_{n+\alpha_f^s}^s,$$

$$F_{n+\alpha_f^f}^f = -\alpha \ddot{u}_{n+\alpha_m^f}^f - 2\xi\omega \dot{u}_{n+\alpha_f^f}^f,$$

where

$$F_{n+\alpha_f^s}^s = \alpha_f^s F_{n+1} + (1 - \alpha_f^s) F_n,$$

$$F_{n+\alpha_f^f}^f = \alpha_f^f F_{n+1} + (1 - \alpha_f^f) F_n.$$

Numerical solution with interpolation

Then

$$\alpha_f^s F_{n+1} + (1 - \alpha_f^s) F_n = (1 - \alpha) \ddot{u}_{n+\alpha_m^s} + \omega^2 u_{n+\alpha_f^s}^s,$$

$$\alpha_f^f F_{n+1} + (1 - \alpha_f^f) F_n = -\alpha \ddot{u}_{n+\alpha_m^f}^f - 2\xi\omega \dot{u}_{n+\alpha_f^f}^f.$$

and

$$(1 - \alpha) [\alpha_m^s \ddot{u}_{n+1}^s + (1 - \alpha_m^s) \ddot{u}_n^s] \\ + \omega^2 [\alpha_f^s u_{n+1}^s + (1 - \alpha_f^s) u_n^s] = \alpha_f^s F_{n+1} + (1 - \alpha_f^s) F_n,$$

$$\alpha_f^f F_{n+1} + (1 - \alpha_f^f) F_n = -\alpha [\alpha_m^f \ddot{u}_{n+1}^f + (1 - \alpha_m^f) \ddot{u}_n^f] \\ - 2\xi\omega [\alpha_f^f \dot{u}_{n+1}^f + (1 - \alpha_f^f) \dot{u}_n^f]$$

$$u_{n+1}^s = u_n^s + \tau \dot{u}_n^s + \tau^2 \left[\left(\frac{1}{2} - \beta^s \right) \ddot{u}_n^s + \beta_s \ddot{u}_{n+1}^s \right],$$

$$\dot{u}_{n+1}^f = \dot{u}_n^f + \tau [\gamma^f \ddot{u}_{n+1}^f + (1 - \gamma^f) \ddot{u}_n^f],$$

$$\dot{u}_{n+1}^s = \dot{u}_n^s + \tau [\gamma^s \ddot{u}_{n+1}^s + (1 - \gamma^s) \ddot{u}_n^s].$$

Numerical solution with interpolation

Finally $\mathbf{v}_{n+1} = A_1^{-1} A_2 \mathbf{v}_n$ with

$$A_1 = \begin{pmatrix} \omega^2 \alpha_f^s & 0 & \frac{(1-\alpha) \alpha_m^s}{\tau^2} & 0 & -\frac{\alpha_f^s}{\tau^2} \\ 0 & 2 \frac{\xi \omega}{\tau} \alpha_f^f & 0 & \frac{\alpha \alpha_m^f}{\tau^2} & \frac{\alpha_f^f}{\tau^2} \\ 1 & 0 & -\beta^s & 0 & 0 \\ 0 & \frac{1}{\tau} & 0 & -\frac{\gamma^f}{\tau} & 0 \\ 0 & \frac{1}{\tau} & -\frac{\gamma^s}{\tau} & 0 & 0 \end{pmatrix},$$

$$A_2 = \begin{pmatrix} -\omega^2(1-\alpha_f^s) & 0 & -\frac{(1-\alpha)(1-\alpha_m^s)}{\tau^2} & 0 & \frac{1-\alpha_f^s}{\tau^2} \\ 0 & -2 \frac{\xi \omega}{\tau} (1-\alpha_f^f) & 0 & \frac{\alpha(\alpha_m^f-1)}{\tau^2} & -\frac{1-\alpha_f^f}{\tau^2} \\ 1 & 1 & \frac{1}{2} - \beta^s & 0 & 0 \\ 0 & \frac{1}{\tau} & 0 & \frac{1-\gamma^f}{\tau} & 0 \\ 0 & \frac{1}{\tau} & \frac{1-\gamma^s}{\tau} & 0 & 0 \end{pmatrix}$$

$$\mathbf{v}_n = (u_n^s, \tau \dot{u}_n^f, \tau^2 \ddot{u}_n^s, \tau^2 \ddot{u}_n^f, \tau^2 F_n)^T.$$

Convergence

- Compute u_{n+1} and expand it in a Taylor series.
- Local error:

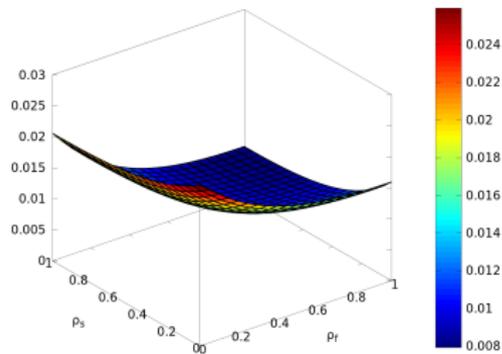
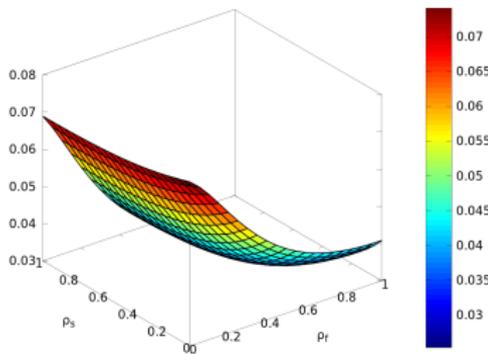
$$\delta_\tau = \tau^2 \frac{\beta^s \gamma^f (\alpha_f^f - \alpha_f^s)}{(1 - \alpha) \alpha_m^s \alpha_f^f \gamma^f + \alpha \alpha_f^s \alpha_m^f \gamma^s} + \mathcal{O}(\tau^3).$$

- If the coefficients of the generalised- α method are chosen with ρ_s and ρ_f the error gets large for the setting

$$\rho_s = 1, \quad \rho_f = 0, \quad \alpha = 0.9.$$

Numerical experiment

- **Setting 1:** $\omega = 1$, $\xi = 0.001$, and $\alpha = 0.5$
- **Setting 2:** $\omega = 1$, $\xi = 0.01$, and $\alpha = 0.8$.
- stepsize $\tau = \pi/15$



Coupling multistep version

First we have

$$\mathbf{u}_{n+1} = \frac{2\alpha_m - 1}{\alpha_m} \mathbf{u}_n - \frac{\alpha_m - 1}{\alpha_m} \mathbf{u}_{n-1} + \frac{\tau(1 - \gamma)}{\alpha_m} \mathbf{f}_{n-1+\alpha_f} + \frac{\tau\gamma}{\alpha_m} \mathbf{f}_{n+\alpha_f}.$$

for the first order problem and

$$\sum_{j=0}^3 [M\alpha_j + \tau C\gamma_j] \mathbf{u}_{n+j} + \tau^2 \sum_{j=0}^2 \delta_j [\mathbf{S}_{n+j+\alpha_f} - \mathbf{F}(t_{n+j+\alpha_f})] = \mathbf{0},$$

for the second order problem.

Coupling multistep version

Then

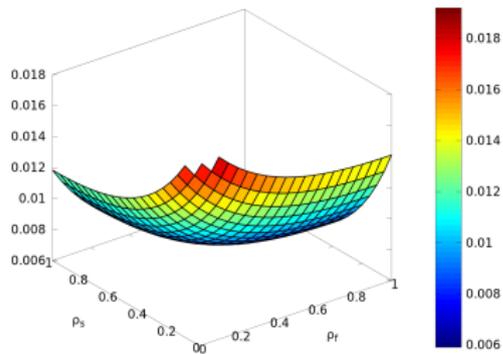
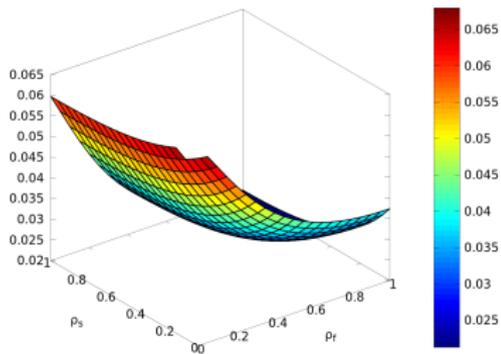
$$(1 - \alpha) \sum_{j=0}^3 \alpha_j^s \mathbf{u}_{n+j}^s + \tau^2 \sum_{j=0}^2 \delta_j [\omega^2 \mathbf{u}_{n+j+\alpha_f} - \mathbf{F}(t_{n+j+\alpha_f})] = 0$$
$$\alpha_m^f \dot{\mathbf{u}}_{n+3}^f - (2\alpha_m^f - 1) \dot{\mathbf{u}}_{n+2}^f + (\alpha_m^f - 1) \dot{\mathbf{u}}_{n+1}^f$$
$$+ \tau(1 - \gamma^f)(2\xi\omega \dot{\mathbf{u}}_{n+1+\alpha_f}^f + \mathbf{f}_{n+1+\alpha_f})$$
$$+ \tau\gamma(2\xi\omega \dot{\mathbf{u}}_{n+2+\alpha_f}^f + \mathbf{f}_{n+2+\alpha_f}) = 0$$
$$\frac{1}{2\tau}(\mathbf{u}_{n+3} - \mathbf{u}_{n+1}) = \dot{\mathbf{u}}_{n+2}$$

Coupling multistep version

It follows

$$\begin{aligned}
 & [(1 - \alpha)\alpha_m^s + \tau^2\beta\omega^2\alpha_f^f]\mathbf{u}_{n+3}^s - \tau^2\beta\omega^2\alpha_f^f\mathbf{F}_{n+3} \\
 &= -[(1 - \alpha)(1 - 3\alpha_m) + \tau^2\left(\frac{1}{2} - 2\beta + \gamma\right)]\omega^2\alpha_f^f + \tau^2\beta\omega^2(1 - \alpha_f^f)]\mathbf{u}_{n+2}^s \\
 & - [(1 - \alpha)(3\alpha_m - 2) + \tau^2\left(\frac{1}{2} + \beta - \gamma\right)]\omega^2\alpha_f^f + \tau^2\left(\frac{1}{2} - 2\beta + \gamma\right)\omega^2(1 - \alpha_f^f)]\mathbf{u}_{n+1}^s \\
 & - [(1 - \alpha)(1 - \alpha_m) + \tau^2\left(\frac{1}{2} + \beta - \gamma\right)]\omega^2(1 - \alpha_f^f)]\mathbf{u}_n^s \\
 & + \tau^2\left[\left(\frac{1}{2} - 2\beta + \gamma\right)\omega^2\alpha_f^f + \beta\omega^2(1 - \alpha_f^f)\right]\mathbf{F}_{n+2} \\
 & + \tau^2\left[\left(\frac{1}{2} + \beta - \gamma\right)\omega^2\alpha_f^f + \tau^2\left(\frac{1}{2} - 2\beta + \gamma\right)\omega^2(1 - \alpha_f^f)\right]\mathbf{F}_{n+1} \\
 & + \tau^2\left(\frac{1}{2} + \beta - \gamma\right)\omega^2(1 - \alpha_f^f)\mathbf{F}_n \\
 & (\alpha_m^f + \tau\gamma 2\xi\omega\alpha_f^f)\dot{\mathbf{u}}_{n+3}^f + \tau\gamma\alpha_f^f\mathbf{F}_{n+3} \\
 &= [(2\alpha_m^f - 1) - \tau\gamma(2\xi\omega(1 - \alpha_f^f) - \tau(1 - \gamma^f)2\xi\omega\alpha_f^f)]\dot{\mathbf{u}}_{n+2}^f \\
 & - [(\alpha_m^f - 1) + \tau(1 - \gamma^f)2\xi\omega(1 - \alpha_f^f)]\dot{\mathbf{u}}_{n+1}^f \\
 & - \tau[\gamma(1 - \alpha_f^f) + (1 - \gamma^f)\alpha_f^f]\mathbf{F}_{n+2} - \tau(1 - \gamma^f)(1 - \alpha_f^f)\mathbf{F}_{n+1} \\
 \dot{\mathbf{u}}_{n+2} &= \frac{1}{2\tau}(\mathbf{u}_{n+3} - \mathbf{u}_{n+1})
 \end{aligned}$$

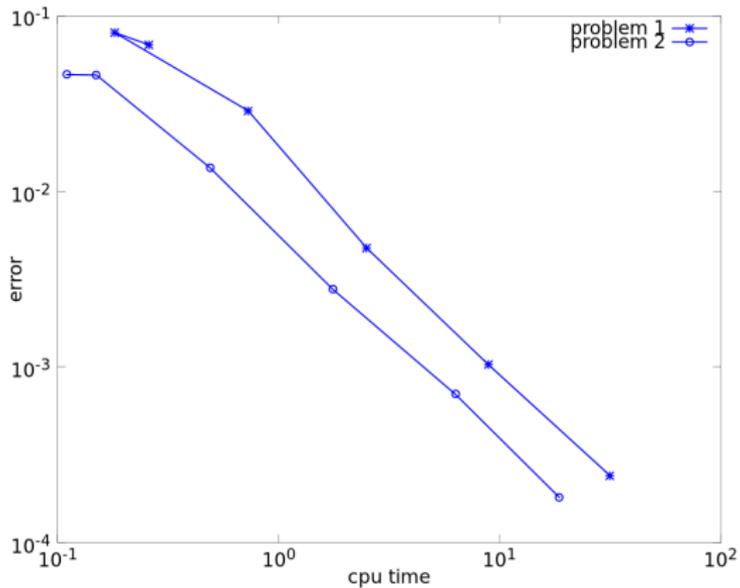
Numerical results



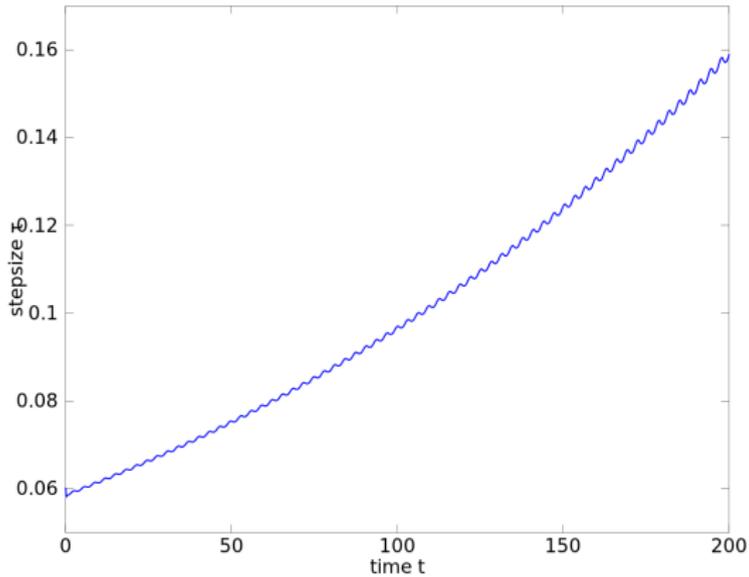
Adaptivity for the partitioned approach

- First compute u_{n+1}^s and \dot{u}_{n+1}^l
- Compute \ddot{u}_{n+1} by evaluating the model problem.
- Compute a second solution with the backward Euler method, i. e. $u_{n+1}^s = u_n^s + \tau \dot{u}_{n+1}^l$ and $\dot{u}_{n+1}^l = \dot{u}_n^l + \tau \ddot{u}_{n+1}$.
- Compute the numerical error r_{n+1} and approximate the new timestep length τ_{n+1} .
- If the numerical error is smaller than the given tolerance the timestep is accepted otherwise it is rejected and has to be recomputed with the new timestep length τ_{n+1} .

Numerical result: CPU versus error



Numerical result: t versus τ



Literature

- J. Chung and G.M. Hulbert: [A time integration algorithm for structural dynamics with improved numerical dissipation: The generalized- \$\alpha\$ method](#). J. Appl. Mech., 60(2):371–375, 1993.
- K. E. Jansen, C. H. Whiting, and G. M. Hulbert: [A generalized- \$\alpha\$ method for integrating the filtered Navier-Stokes equations with a stabilized finite element method](#). Comput. Methods Appl. Mech. Eng., 190(3-4):305–319, 2000.
- M.M. Joosten, W.G. Dettmer, and D. Perić: [On the temporal stability and accuracy of coupled problems with reference to fluid–structure interaction](#). Int. J. Numer. Methods Fluids, 64(10-12):1363–1378, 2010.