



Technische Universität Braunschweig



Introduction to Scientific Computing

Overview

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What we have learned?

Real world







Real world \Rightarrow ODEs



Felix's free fall

$$m\frac{d^2y}{dt^2} = \frac{1}{2}\rho C_{\rm D} A \left(\frac{dy}{dt}\right)^2 - mg$$

where *m* is mass, C_D drag coefficient, ρ air density and *g* is gravitational constant

Here, the maximal order of derivative is 2, so we say that the ODE is of the **second** order.





Higher order ODE \Rightarrow first order ODE

Higher order ODEs can be solved by using the method of transformation to the system of lower order ODEs.

$$m\frac{d^2y}{dt^2} = \frac{1}{2}\rho C_{\rm D} A \left(\frac{dy}{dt}\right)^2 - mg$$

Take $x_1 := y$, $x_2 = \frac{dy}{dt}$ as a new variable, then

$$\frac{dx_1}{dt} = \frac{dy}{dt}, \quad \frac{dx_2}{dt} = \frac{d^2y}{dt^2}$$

which further leads to the system of first order ODEs:

$$\dot{\mathbf{x}} = \begin{pmatrix} \frac{dx_1}{dt} \\ \frac{dx_2}{dt} \end{pmatrix} = \begin{pmatrix} x_2 \\ \frac{1}{2m} \rho C_{\mathrm{D}} A(x_1)^2 - g \end{pmatrix} = \mathbf{f}(\mathbf{x})$$





First order ODE \Rightarrow stability

Equilibria points are obtained by solving

 $\bm{f}(\bm{x}_*) = \bm{0}$

which possibly requires Newton method:

$$\mathbf{x}_{*}^{(k+1)} = \mathbf{x}_{*}^{(k)} - (\mathbf{f}'(\mathbf{x}^{(k)}))^{-1}\mathbf{f}(\mathbf{x}^{(k)})$$

Stability is determined by

 $\dot{\boldsymbol{x}} = \boldsymbol{\textit{D}} \boldsymbol{f}(\boldsymbol{x}_*) \boldsymbol{x}$

where $\mathbf{J} := D\mathbf{f}(\mathbf{x}_*)$ is the Jacobi matrix.





ODE stability

$$\mathbf{J} = \mathbf{Q} e^{t \wedge} \mathbf{Q}^{-1}$$

Thus,

- if any $\operatorname{Re}(\lambda_j) > 0$ then $x \to \infty$ when $t \to \infty$ (unstable)
- if for almost all Re (λ_j) < 0 besides at least one Re (λ_j) = 0 then it is stable but not attractive in case of linear system, otherwise it is difficult to say based only on the first order Taylor expansion
- if **all** Re $(\lambda_j) < 0$ then the state x_* is stable and attractive

So find eigenvalues of J and check if you are on the left side of complex plane!!





where we are







Algebraic solving is possible for linear system of ODEs. Solution is of the form

$$\mathbf{x} = c_1 e^{\lambda t} \mathbf{v} + c_2 (t e^{\lambda t} \mathbf{v} + e^{\lambda t} \mathbf{w})$$

in which \mathbf{v} is eigenvector and \mathbf{w} is the generalised eigenvector. These are obtained by solving

$$(A - \lambda I)\mathbf{v} = \mathbf{0}, \quad (A - \lambda I)\mathbf{w} = \mathbf{v}.$$

However, the algebraic solution is not something that is used in practice.





First order ODE \Rightarrow numerical integration

$$\dot{x} = f(t) \Rightarrow x(t_n) = x(t_m) + \int_{t_m}^{t_n} f(t) dt$$

and thus includes the numerical integration

$$\int_{t_m}^{t_n} f(t) dt = \int_{t_m}^{t_n} P_n(t) dt = \sum_{i=1}^N w_i f(t_i)$$

in which P_n is *n*-th order interpolation polynomial satisfying interpolation conditions

$$P_n(t_i) = f(t_i), \quad i = 1, ..., n+1$$

The interpolation points t_i are chosen to be equidistant or are assumed to be unknown (in this case they are computed from the error

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$\textbf{ODE} \Rightarrow \textbf{integration} \Rightarrow \textbf{difference equation}$

Solution

$$\mathbf{x}(t_n) = \mathbf{x}(t_m) + \int_{t_m}^{t_n} f(t) dt$$

is difference equation. If f(t) is approximated by constant, then one may distinguish explicit from implicit Euler method:

$$x(t_{n+1}) = x(t_n) + hf(x_n, t_n)$$
 and $x(t_{n+1}) = x(t_n) + hf(x_{n+1}, t_{n+1})$

These are one step methods described by interpolation points $t = t_n$ or $t = t_{n+1}$. Multistep methods use approximation of *n*-th order to model the right hand side $f \approx P_n$ or the state itself $x \approx P_n$, $\frac{dP_n}{dt} = f$. Hence, they need several interpolation points and thus several starting points.





Numerical scheme \Rightarrow convergence

Convergence implies consistency and zero stability. Consistency does not imply convergence.

The method is said to be consistent if the local (truncation) error converges to zero when the step-size becomes arbitrarily small, i.e.

$$\lim_{h\to 0} T_h = \lim_{h\to 0} \frac{\epsilon_{loc}}{h} = 0$$

The scheme is consistent of order *p* if

$$\max \| \frac{\epsilon_{\textit{loc}}}{h} \| \leqslant Ch^{p}$$

The local error ϵ_{loc} is the difference between the exact $x_a(t_{n+1})$ and numerical $x(t_{n+1})$ solutions at t_{n+1} where $x(t_{n+1})$ is computed by a numerical scheme starting from $x_a(t_n)$.

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Linear multistep method is said to be convergent if

$$\max_{t_n \in T} \|\epsilon_{glob}(t_n, h)\| \to 0, \quad h \to 0$$

and it has the order of convergence q if

$$\max_{t_n \in T} \|\epsilon_{glob}(t_n, h)\| \leqslant Ch^q$$

where the global error is $\epsilon_{glob} = x_a(t) - x(t)$ in which $x_a(t)$ is the exact solution of the differential equation $\dot{x} = f(t, x)$ and x(t) is the approximate solution. *C* is a constant independent of *h*





Numerical scheme \Rightarrow step size

The difference equation $\sum_{j=0}^{k} a_j x_{n+j} = h \sum_{j=0}^{k} b_j f(t_{n+j}, x_{n+j})$ is stable if roots

 $\rho(\xi) - z\sigma(\xi) = 0$

are smaller by amplitude than 1, or eventually only one of them is equal to 1. With respect to this one defines

Definition

The region of absolute stability for the LMM is the set of points $z = h\lambda$ in the complex plane for which the polynomial $\rho(\xi) - z\sigma(\xi) = 0$ satisfies the root condition.

$$G_{s} := \{ z \in \mathbb{C} \ : \ |\xi(z)| \leqslant 1 \}$$





where we are

Real world







Numerical scheme \Rightarrow solving system

Applying implicit numerical scheme onto ODE

$$\sum_{j=0}^{k} a_{j} x_{n+j} = h \sum_{j=0}^{k} b_{j} f(t_{n+j}, x_{n+j})$$

such as for example implicit Euler method

$$x_{n+1} = x_n + hf(x_{n+1}, t_{n+1})$$

one ends up with the linear/nonlinear system of equations

$$g(x_{n+1}) = x_n + hf(x_{n+1}, t_{n+1}) - x_{n+1} = 0$$

in which x_{n+1} is unknown.

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Solving system \Rightarrow fixed point iteration

The solution of

$$g(z) = 0$$

is fixed point z_* in interval $[z_a, z_b]$ which satisfies

$$z_* = \varphi(z_*)$$

The solution is unique and exists if **Banach fixed point** theorem is valid, i.e.

- $\varphi(z): [z_a, z_b] \rightarrow [z_a, z_b]$
- the mapping φ(z) is Lipshitz continuous

$$\|\varphi(z) - \varphi(\hat{z})\| \leqslant q \|z - \hat{z}\|$$

the mapping is contractive q < 0





Solving system \Rightarrow fixed point iteration

We have studied two error estimates

A posteriori error estimate:

$$\|\mathbf{z}_* - \mathbf{z}_k\| \leqslant rac{q}{1-q} \|\mathbf{z}_k - \mathbf{z}_{k-1}\|$$

A priori error estimate:

$$\|\mathbf{z}_* - \mathbf{z}_k\| \leqslant \frac{q^k}{1-q} \|\mathbf{z}_1 - \mathbf{z}_0\|.$$

Here, index k is not time step, but iteration number.





Solving system \Rightarrow linear system

System:

$$\mathbf{A}\mathbf{z} = \mathbf{b}$$

Gauss-Seidel method:

$$\mathbf{Lz}^{(k+1)} = \mathbf{b} - \mathbf{Uz}^{(k)}$$

Jacobi method:

$$\mathbf{D}\mathbf{z}^{(k+1)} = \mathbf{b} - \mathbf{R}\mathbf{z}^{(k)}$$

Successive over-relaxation (SOR) method:

$$\mathbf{z}^{(k+1)} = (\mathbf{D} + \omega \mathbf{L})^{-1} \big(\omega \mathbf{b} - [\omega \mathbf{U} + (\omega - 1)\mathbf{D}]\mathbf{z}^{(k)} \big).$$

Conjugate gradients

$$\mathbf{z}^{(k+1)} = \mathbf{z}^{(k)} + \alpha^{(k)} \mathbf{d}^{(k)}$$





Solving system \Rightarrow nonlinear system

$$F(z) = 0$$

 $F(z) = F(z_0) + F'(z_0)(x - z_0) + h.o.t. = 0$

Jacobian:

$$\mathbf{J}(\mathbf{z}) = \mathbf{F'}(\mathbf{z})$$

with elements:

$$J_{ij}=\frac{\partial f_i}{\partial z_j}$$

Newton iteration:

$$\mathbf{z}_k = \mathbf{z}_{k-1} - \frac{\mathbf{F}(\mathbf{z}_{k-1})}{\mathbf{J}(z_{k-1})}$$





Solving system \Rightarrow nonlinear system







Solving system \Rightarrow nonlinear system

Besides classical Newton algorithm one may also use:

simplified (stationary) Newton

$$\mathbf{z}_k = \mathbf{z}_{k-1} - \frac{\mathbf{F}(\mathbf{z}_{k-1})}{\mathbf{J}(z_0)}$$

modified Newton

$$\mathbf{z}_k = \mathbf{z}_{k-1} - m \frac{\mathbf{F}(\mathbf{z}_{k-1})}{\mathbf{J}(\mathbf{z}_{k-1})}$$

Newton method with restart

$$\mathbf{z}_k = \mathbf{z}_{k-1} - \frac{\mathbf{F}(\mathbf{z}_{k-1})}{\mathbf{J}(z_m)}$$

Broyden's method (updating Jacobian)





Modelling, as well as numerical solving need to satisfy the property of well-posedness:

- solution exists
- solution is unique
- small perturbations in data cause small pertubrations in the solution





Solving system \Rightarrow condition number

Model

$$y = f(x)$$

has condition number

$$k(x) = \left| \frac{\partial f}{\partial x} \frac{x}{f(x)} \right|$$

The system

$$Ax = b$$

has condition number

$$k(A) = \|A\| \|A^{-1}\|$$





where we are

Real world







DIFFERENCE AND DIFFERENTIAL EQUATIONS







Difference vs. Differential equations

Discrete system	Continuous system
Difference	Differential
$x_{n+1}-x_n=2h^2n$	$\frac{dx}{dt} = 2t$
	Discretise in time
	$x_{n+1} - x_n = 2h^2n$

Differential equations become difference when we apply numerical integration (time discretisation)!





Homogeneous equations (zero rhs)

Difference	Differential
$x_{n+1}-x_n=0$	$\frac{dx}{dt} = 0$
Solution form	Solution form
$x_n = C \rho^n$	$x(t) = Ce^{ ho t}$
Characteristic	Characteristic
$\rho-1=0$	ho=0
$x_n = C(1)^n = C$	$x = Ce^{0t} = C$

I.





Homogeneous solution wrt roots

Difference	Differential
Single roots p _i	Single roots ρ_i
$x_n = \sum_i C_i \rho_i^n$	$x(t) = \sum_{i} C_{i} e^{\rho_{i} t}$
Double roots ρ_1	Double roots ρ_1
$x_n = C_1 \rho_1^n + n C_2 \rho_1^n$	$x(t) = C_1 e^{\rho_1 t} + t C_2 e^{\rho_1 t}$

Roots and eigenvalues are equal!!





Homogeneous solution wrt roots

Difference	Differential
Complex $\rho = \alpha \pm \beta i$	Complex $\rho = \alpha \pm \beta i$
$\mu := \sqrt{\alpha^2 + \beta^2}$	$e^{(\alpha\pm\beta i)t} = e^{\alpha t}e^{\beta it}$
$\tan\varphi=\beta/\alpha$	
$x_n = \mu^n (\mathbf{A}\cos(n\phi) + i\mathbf{B}\sin(n\phi))$	$\mathbf{x} = \mathbf{e}^{\alpha t} (\mathbf{A} \cos(\beta t) + i\mathbf{B} \sin(\beta t))$
Note that in one case is the modulus and in another real part of the	



Non-homogeneous- Particular solution



Differential

$$x = x_h + x_p$$

 $x_h = Ce^{pt}$,

variation of constant

$$x_p = c(t)e^{pt}$$





STABILITY







Stability

Difference	Differential
$x_{n+1} = F(x_n)$	$\frac{dx}{dt} = f(x)$
$x_* = F(x_*)$	$f(x_*)=0$
roots of charact. equation or	roots of charact. equation or
$[\lambda, v] = ext{eig } DF(x_*)$	$[\lambda, v] = ext{eig } Df(x_*)$
$(\lambda)^n \Rightarrow \mathrm{abs}\ (\lambda) \leqslant 1$ (real)	${oldsymbol e}^{\lambda t} \Rightarrow (\lambda) \leqslant {f 0}$ (real)
$abs (\mu) \leqslant 1$ (complex)	$(\textit{Re}(\lambda)) \leqslant 0$ (complex)





In this lecture we have only scratched the surface of the process of solving ordinary differential equations and stability properties of numerical schemes. Next semester will deal more seriously with the process of time discretisation of ordinary differential equations and time dependent partial differential equations. Special attention is put on the so-called stiff-systems which are very often arising in practical applications, as well as differential algebraic equations (differential equations with algebraic conditions).





for those interested in applied mathematics there are possiblities of offering student project (3months) or master (6 months) topics. The focus are ODEs/PDEs which describe

- plasticity, viscoplasticity, damage (description of cellulose, concrete, metals etc.)
- weather prediction models (different Lorenz models)
- simple biological models (modelling insulin change in body, drug transport through blood)

Topics are: uncertainty quantification, identification, control, adaptive algorithms, etc.





Thats all folks



Hints can never hurt you, except on exams. Then, I can mess with your mind.



