

Fourier series

Projections in $L_2([0, 2\bar{u}])$ \leftarrow infinitely dimensional function space
rep. projection in \mathbb{R}^d \leftarrow d-dimensional Euclidean space

$$\vec{x} = \begin{pmatrix} x_1 \\ x_2 \\ \vdots \\ x_d \end{pmatrix} = x_1 \begin{pmatrix} 1 \\ 0 \\ \vdots \\ 0 \end{pmatrix} + x_2 \begin{pmatrix} 0 \\ 1 \\ \vdots \\ 0 \end{pmatrix} + \dots + x_d \begin{pmatrix} 0 \\ 0 \\ \vdots \\ 1 \end{pmatrix} = \sum_{k=1}^d x_k \vec{e}_k$$

unit vectors

$\vec{P}_k = x_k \vec{e}_k$ is the projection of \vec{x} onto \vec{e}_k

$\lambda_k?$ $\vec{e}_k \perp \vec{x} - \lambda_k \vec{e}_k$

$$\vec{e}_k \cdot (\vec{x} - \lambda_k \vec{e}_k) = 0 \iff \vec{e}_k \cdot \vec{x} = \lambda_k \vec{e}_k \cdot \vec{e}_k$$

$$\vec{P}_k = \lambda_k \vec{e}_k = \frac{\langle \vec{x}, \vec{e}_k \rangle_{\mathbb{R}^d}}{\|\vec{e}_k\|_2^2} \vec{e}_k = x_k \vec{e}_k$$

again in a more general form : Let $\{\vec{v}_1, \dots, \vec{v}_d\}$ be an orthonormal system

linear combination $\vec{x} = \lambda_1 \vec{v}_1 + \dots + \lambda_d \vec{v}_d$? $\lambda_1, \dots, \lambda_d$

multiply by $\dots \vec{v}_k$: $\langle \vec{x}, \vec{v}_k \rangle_{\mathbb{R}^d} = \lambda_1 \langle \vec{v}_1, \vec{v}_k \rangle_{\mathbb{R}^d} + \dots + \lambda_d \langle \vec{v}_d, \vec{v}_k \rangle_{\mathbb{R}^d}$

$$= \lambda_1 \delta_{1,k} + \dots + \lambda_d \delta_{d,k} = \lambda_k$$

with Kronecker symbol

$$\delta_{k,l} = \begin{cases} 0 & \text{if } k \neq l \\ 1 & \text{if } k = l \end{cases}$$

i.e. $\langle \vec{x}, \vec{v}_k \rangle_{\mathbb{R}^d} = \lambda_k$

$$\vec{x} = \langle \vec{x}, \vec{v}_1 \rangle_{\mathbb{R}^d} \vec{v}_1 + \dots + \langle \vec{x}, \vec{v}_d \rangle_{\mathbb{R}^d} \vec{v}_d$$

$$\vec{x} = \sum_{k=1}^d \underbrace{\langle \vec{x}, \vec{v}_k \rangle_{\mathbb{R}^d}}_{\text{projection of } \vec{x} \text{ onto } \vec{v}_k} \vec{v}_k$$

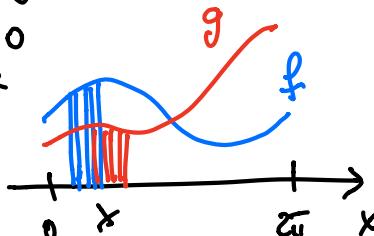
energy

rep. Lebesgue function space $L_2([0, 2\bar{u}])$ contains all "functions" $f: [0, 2\bar{u}] \rightarrow \mathbb{R}$ or \mathbb{C} with

$$\int_0^{2\bar{u}} |f(x)|^2 dx < \infty$$

Scalar product $\langle f, g \rangle_{L_2([0, 2\bar{u}])} = \int_0^{2\bar{u}} f(x) \overline{g(x)} dx$

norm $\|f\|_{L_2([0, 2\bar{u}])} = \sqrt{\langle f, f \rangle_{L_2([0, 2\bar{u}])}}$



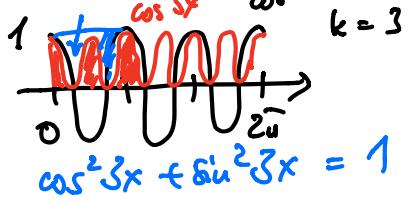
$\mathbb{C}^d: \langle \vec{a}, \vec{b} \rangle = \sum_{k=1}^d a_k \overline{b_k}$

$\vec{a} = \begin{pmatrix} a_1 \\ a_2 \\ \vdots \\ a_d \end{pmatrix}$, $\vec{b} = \begin{pmatrix} b_1 \\ b_2 \\ \vdots \\ b_d \end{pmatrix}$

$$\langle \vec{a}, \vec{b} \rangle = a_1 b_1 + a_2 b_2 + \dots + a_d b_d$$

$$= \left(\int_0^{2\pi} |f(x)|^2 dx \right)^{1/2} = \sqrt{\pi}$$

Ex. $f(x) = \cos kx$, $k \in \mathbb{R}$, $k \neq 0$, $\|f\|_{L_2([0, 2\pi])} = \left(\int_0^{2\pi} \cos^2 kx dx \right)^{1/2} = \sqrt{\frac{2}{\pi}}$



projection of $g \in \mathbb{R}$ onto $f(x) = \cos kx$

$$\lambda_f(x) = \frac{\langle g, f \rangle_{L_2([0, 2\pi])}}{\|f\|_{L_2([0, 2\pi])}^2} f(x)$$

$$= \frac{1}{\pi} \int_0^{2\pi} g(x) \cos kx dx \cdot \cos kx$$

λ coefficient how strongly $f(x) = \cos kx$ is contained in g .

Real Fourier series

Joseph Fourier (1768 Auxerre - 1830 Paris)

Theorem: The functions $\left\{ \frac{1}{\sqrt{2\pi}}, \frac{1}{\sqrt{\pi}} \cos kx, \frac{1}{\sqrt{\pi}} \sin kx : k=1, 2, \dots \right\}$ form an orthonormal system in $L_2([0, 2\pi])$.

Proof: $\langle \frac{1}{\sqrt{\pi}} \cos kx, \frac{1}{\sqrt{\pi}} \cos lx \rangle_{L_2([0, 2\pi])} = \frac{1}{\pi} \int_0^{2\pi} \cos kx \cdot \cos lx dx = \frac{1}{2\pi} \int_0^{2\pi} \underbrace{\cos(kx+lx)}_{\cos kx \cos lx - \sin kx \sin lx} + \underbrace{\cos(kx-lx)}_{\cos kx \cos lx + \sin kx \sin lx} dx \rightarrow \cos kx \cos l + \sin kx \sin l$

$$= \frac{1}{2\pi} \int_0^{2\pi} \cos((k+l)x) dx + \frac{1}{2\pi} \int_0^{2\pi} \cos((k-l)x) dx = \delta_{k,l}$$

$= 0$ due to $k, l \geq 1$ $= 0$ for $k-l \neq 0$
 $= 2\pi$ for $k-l=0$

and $\langle \frac{1}{\sqrt{\pi}} \cos kx, \frac{1}{\sqrt{2\pi}} \rangle_{L_2([0, 2\pi])} = \frac{1}{\pi\sqrt{2}} \int_0^{2\pi} \cos kx dx = 0$

and etc. for all other combinations. \square

Theorem: The functions $\left\{ \frac{1}{\sqrt{2\pi}}, \frac{1}{\sqrt{\pi}} \cos kx, \frac{1}{\sqrt{\pi}} \sin kx : k=1, 2, \dots \right\}$ span a linear hull, which is dense in $L_2([0, 2\pi])$. (without proof)

i.e. every function $f \in L_2([0, 2\pi])$ can be approximated by a function from the linear hull with arbitrary accuracy.

$$f(x) = \underbrace{\left\langle f, \frac{1}{\sqrt{2\pi}} \right\rangle}_{a_0/2} \frac{1}{\sqrt{2\pi}} + \sum_{k=1}^{\infty} \underbrace{\left\langle f, \frac{1}{\sqrt{\pi}} \cos kx \right\rangle}_{a_k} \frac{1}{\sqrt{\pi}} \cos kx + \sum_{k=1}^{\infty} \underbrace{\left\langle f, \frac{1}{\sqrt{\pi}} \sin kx \right\rangle}_{b_k} \frac{1}{\sqrt{\pi}} \sin kx$$

$$f(x) = \frac{a_0}{2} + \sum_{k=1}^{\infty} a_k \cos kx + b_k \sin kx //$$

Fourier Series

$k=1$

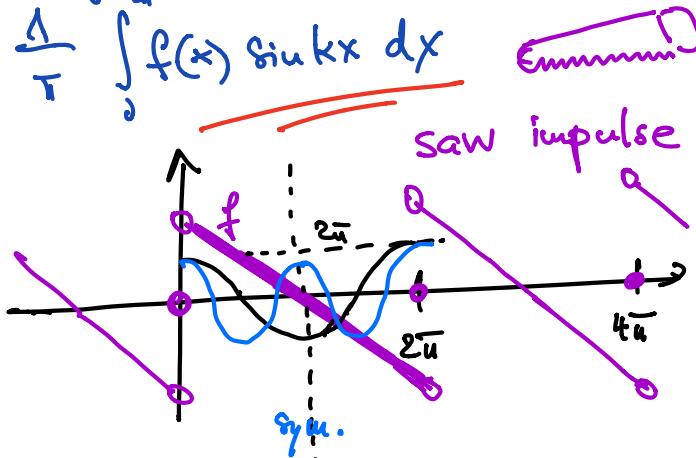
$$\text{Fourier coefficients } \frac{a_0}{2} = \left\langle f, \frac{1}{\sqrt{2\pi}} \right\rangle_{L^2} \frac{1}{\sqrt{2\pi}} = \frac{1}{2\pi} \int_0^{2\pi} f(x) dx, \quad a_0 = \frac{1}{\pi} \int_0^\pi f(x) dx \quad (*)$$

and $a_k = \left\langle f, \frac{1}{\sqrt{\pi}} \cos kx \right\rangle_{L^2} \frac{1}{\sqrt{\pi}} = \frac{1}{\pi} \int_0^\pi f(x) \cos kx dx$

$b_k = \left\langle f, \frac{1}{\sqrt{\pi}} \sin kx \right\rangle_{L^2} \frac{1}{\sqrt{\pi}} = \frac{1}{\pi} \int_0^\pi f(x) \sin kx dx$ Eeeeeeee

Exp.

$$f(x) = \begin{cases} 0 & \text{for } x = 2\bar{u} \\ \frac{1}{2}(\bar{u}-x) & \text{for } x \in (0, 2\bar{u}) \\ 2\bar{u} - \text{periodically expanded} & \end{cases}$$



$$a_k = \frac{1}{\pi} \int_0^\pi \frac{1}{2}(\bar{u}-x) \cos kx dx$$

$$= \frac{1}{\pi} \int_0^\pi \frac{1}{2}(\bar{u}-x) \cos kx dx + \underbrace{\frac{1}{\pi} \int_\pi^{2\pi} \frac{1}{2}(\bar{u}-x) \cos kx dx}_{0}$$

$$= \frac{1}{\pi} \int_0^{2\bar{u}} \frac{\bar{u}-x}{2} \cos kx dx - \frac{1}{\pi} \int_\pi^{2\bar{u}} \frac{x-\bar{u}}{2} \cos(k(2\bar{u}-x)) dx \quad \text{with } \frac{dx}{dz} = -1$$

Substitution $x = 2\bar{u}-z$ $\cos kx$

$$b_k = \frac{1}{\pi} \int_0^{2\bar{u}} \frac{\bar{u}-x}{2} \sin kx dx \quad \text{i.byparts}$$

with $u' = -1$
 $v = -\frac{\cos kx}{k}$

$$= \frac{1}{\pi} \int_0^{2\bar{u}} \frac{\bar{u}-x}{2} \sin kx dx = \frac{x-\bar{u}}{2\pi} \frac{\cos kx}{k} \Big|_{x=0}^{2\bar{u}} - \frac{1}{2\pi} \int_0^{2\bar{u}} \frac{\cos kx}{k} dx$$

$$= \frac{1}{2\pi} (2\bar{u}-\bar{u}) \frac{\cos 2\bar{u}k}{k} - \frac{1}{2\pi} (0-\bar{u}) \frac{\cos 0}{k} \quad u'v = 0$$

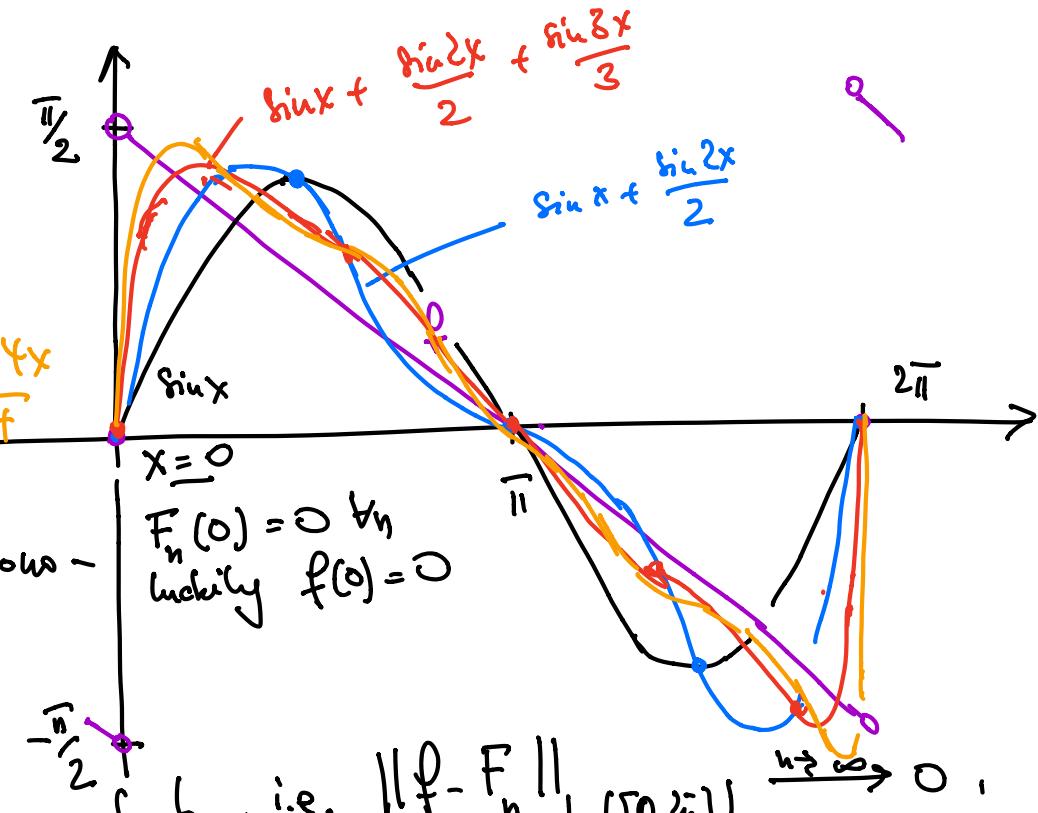
$$= \frac{1}{k} \frac{1}{2} \quad \text{i.e. } b_k = \frac{1}{k}$$

$$\Rightarrow f(x) = \sum_{k=1}^{\infty} b_k \sin kx = \sum_{k=1}^{\infty} \frac{\sin kx}{k} = \sin x + \frac{\sin 2x}{2} + \frac{\sin 3x}{3} + \dots$$

cf. divergent harmonic series $1 + \frac{1}{2} + \frac{1}{3} + \frac{1}{4} + \dots \rightarrow \infty$

partial sums

- $\sin x$
- $\sin x + \frac{\sin 2x}{2}$
- $\sin x + \frac{\sin 2x}{2} + \frac{\sin 3x}{3}$
- \dots
- $\sin x + \frac{\sin 2x}{2} + \frac{\sin 3x}{3} + \frac{\sin 4x}{4}$



Fourier polynomials (in particular polynomials)

$$F_n(x) = \sum_{k=1}^{n+1} \frac{\sin kx}{k}$$

converge to f in the sense of

$$\|f - F_n\|_{L_2([0, 2\pi])} \xrightarrow{n \rightarrow \infty} 0,$$

Complex Fourier Series

$$f: [0, 2\pi] \rightarrow \mathbb{C}$$

Theorem: The functions $\left\{ \frac{1}{\sqrt{2\pi}} e^{ikx}, k \in \mathbb{Z} \right\}$ form an orthonormal system.

- wikipedia : concepts
- scripts: calculus 1 & 2
- textbooks on calculus
e.g. tu-bs.de → library
→ Springer Link **basic idea**
- So einfach ist Mathematik ...
ODE / PDE

Proof: $\left\langle \frac{1}{\sqrt{2\pi}} e^{ikx}, \frac{1}{\sqrt{2\pi}} e^{ilx} \right\rangle_{L_2([0, 2\pi])} =$
 $\int_0^{2\pi} \frac{1}{\sqrt{2\pi}} e^{ikx} \cdot \frac{1}{\sqrt{2\pi}} e^{ilx} dx = \int_0^{2\pi} \frac{1}{2\pi} e^{ikx} e^{-ilx} dx$
 $= \frac{1}{2\pi} \int_0^{2\pi} e^{i(k-l)x} dx = \begin{cases} 0 & \text{for } k \neq l \\ \frac{1}{2\pi} \int_0^{2\pi} 1 dx = 1 & \text{for } k = l \end{cases}$
 $= \delta_{k,l}$

again projection \square

$$f(x) = \sum_{k \in \mathbb{Z}} \underbrace{\left\langle f, \frac{1}{\sqrt{2\pi}} e^{ikx} \right\rangle}_{\text{functions in the "base"}} \cdot \underbrace{\frac{1}{\sqrt{2\pi}} e^{ikx}}_{c_k}$$
$$= \sum_{k=-\infty}^{\infty} \underbrace{\frac{1}{2\pi} \left\langle f, e^{ikx} \right\rangle}_{\tilde{c}_k} e^{ikx}$$

conjugate complex number

Complex Fourier coefficient $c_k = \frac{1}{2\pi} \int_0^{2\pi} f(x) e^{-ikx} dx$

Fourier coefficients

$$f(x) = \sum_{k \in \mathbb{Z}} c_k e^{ikx} = c_0 + \sum_{k=1}^{\infty} (c_k e^{ikx} + c_{-k} e^{-ikx})$$

compare $\dots = +c_{-2} e^{-2ix} + c_{-1} e^{-ix} + c_0 + c_1 e^{ix} + c_2 e^{2ix} + \dots$
 $\dots = c_0 + (c_1 e^{ix} + c_{-1} e^{-ix}) + (c_2 e^{2ix} + c_{-2} e^{-2ix}) + \dots$

$$= c_0 + \sum_{k=1}^{\infty} c_k (\cos kx + i \sin kx) + c_{-k} (\cos kx - i \sin kx)$$

$$= c_0 + \underbrace{\sum_{k=1}^{\infty} (c_k + c_{-k})}_{\frac{a_0}{2}} \cos kx + \underbrace{i(c_k - c_{-k})}_{b_k} \sin kx$$

$$\Rightarrow \frac{a_0}{2} = c_0, \quad a_k = c_k + c_{-k}, \quad b_k = i(c_k - c_{-k}) \quad \text{for } k=1, 2, \dots$$

and $c_k = \frac{1}{2}(a_k - i b_k), \quad c_{-k} = \frac{1}{2}(a_k + i b_k) \quad \text{for } k=1, 2, \dots$

Ex. $f(x) = \begin{cases} 0 & \text{for } x = 2\pi \\ \frac{1}{2}(\pi - x) & \text{for } x \in (0, 2\pi) \\ \text{2}\pi\text{-periodically extended} & \end{cases}$ with $a_k = 0, b_k = \frac{1}{k}$

$$\Rightarrow c_0 = 0, \quad c_k = \frac{-i}{2k}, \quad c_{-k} = \frac{i}{2k} = \frac{-i}{2(-k)}, \quad \tilde{f}(x) = \sum_{k=-\infty, k \neq 0}^{\infty} \frac{-i}{2k} e^{ikx}$$

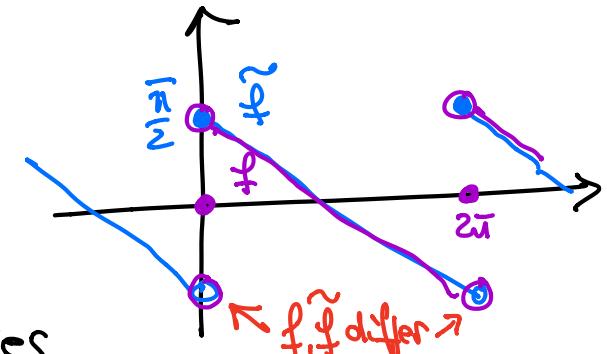
Remark $\tilde{f}(x) = \begin{cases} \frac{\pi-x}{2} & \text{for } x \in [0, 2\pi] \\ \text{2}\pi\text{-periodically extended} & \end{cases}$

$$\tilde{f}(0) = 0, \quad \tilde{f}(0) = \frac{\pi}{2}$$

$$\tilde{f}(x) = \sum_{k=1}^{\infty} \frac{\sin kx}{k} \quad \text{same Fourier series}$$

not pointwise but in the sense of L_2 , i.e.

$$f \stackrel{L_2}{=} g \iff \|f - g\|_{L_2([0, 2\pi])} = \left(\int_0^{2\pi} |f(x) - g(x)|^2 dx \right)^{1/2} = 0$$



single points do
not affect the
integral

Fourier series converge in the sense of L_2 , typically not pointwise!

derivatives, $f \in C^1([0, 2\pi])$, i.e. once differentiable function f

$$f(x) = \sum_{k=-\infty}^{\infty} c_k e^{ikx} = \frac{a_0}{2} + \sum_{k=1}^{\infty} a_k \cos kx + b_k \sin kx \quad (a_k \xrightarrow{k \rightarrow \infty} 0, b_k \xrightarrow{k \rightarrow \infty} 0)$$

$$f'(x) = \sum_{k=-\infty}^{\infty} ik c_k e^{ikx} = \sum_{k=1}^{\infty} -k a_k \sin kx + k b_k \cos kx \quad (ka_k \xrightarrow{k \rightarrow \infty} 0, kb_k \xrightarrow{k \rightarrow \infty} 0)$$

The smoother the function f is,
the faster decay the Fourier coefficients.

Frequencies

regard only one summand of Fourier series $g(x) = a_k \cos kx + b_k \sin kx$
harmonic oscillation, $a_k, b_k \in \mathbb{R}$ \xrightarrow{k} frequency \xrightarrow{k}

ansatz $g(x) = A_k \cos(kx - \varphi_k)$, amplitude A_k , phase φ_k

$$A_k \cos(kx - \varphi_k) = \underbrace{A_k \cos \varphi_k}_{a_k} \cos kx + \underbrace{A_k \sin \varphi_k}_{b_k} \sin kx$$

$$a_k = A_k \cos \varphi_k, b_k = A_k \sin \varphi_k$$

$$A_k = \sqrt{a_k^2 + b_k^2} = \sqrt{(A_k \cos \varphi_k)^2 + (A_k \sin \varphi_k)^2} = A_k \sqrt{\cos^2 \varphi_k + \sin^2 \varphi_k}$$

$$\tan \varphi_k = \frac{\sin \varphi_k}{\cos \varphi_k} = \frac{b_k}{a_k} \quad \text{or} \quad \varphi_k = \arctan \frac{b_k}{a_k}$$

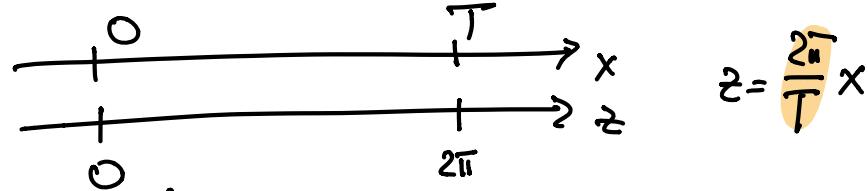
More general periodic functions

g might be a function of period T , $g(x) = g(x+T) \quad \forall x$

Search for $g(x) = \frac{a_0}{2} + \sum_{k=1}^{\infty} a_k \cos \frac{2\pi k x}{T} + b_k \sin \frac{2\pi k x}{T}$

coordinate transformation

$$x=0 \mapsto z=0, x=T \mapsto z=2\pi$$



$f(z) = g(x)$ is $\frac{2\pi}{T}$ -periodic function

Fourier coefficients of f : $\tilde{a}_k(f) = \frac{1}{T} \int f(z) \cos kz dz$

Substitution $x = \frac{T}{2\pi} z$, $dx = \frac{T}{2\pi} dz$

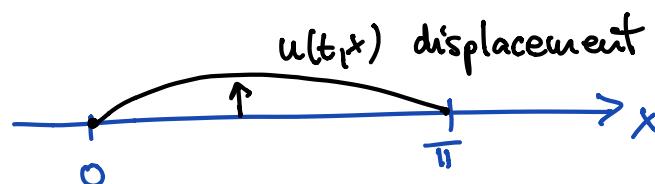
$$a_k(g) = \frac{1}{\pi} \int_0^T g(x) \cos \frac{2\pi k x}{T} \cdot \frac{2\pi}{T} dx$$

$$= \frac{2}{T} \int_0^T g(x) \cos \frac{2\pi k x}{T} dx$$

analogously $b_k(g) = \frac{2}{T} \int_0^T g(x) \sin \frac{2\pi k x}{T} dx$

Fourier series in eigen-oscillations

String



$$\rho \frac{\partial^2}{\partial t^2} u(t, x) = P \frac{\partial^2}{\partial x^2} u(t, x)$$

↑ density ↑ acceleration ↑ tension

boundary condition $u(t, 0) = u(t, \pi) = 0$

look for harmonic oscillations (eigen-oscillations) $u(t, x) = \cos \omega t \cdot V(x)$

$$\frac{\partial^2}{\partial t^2} u = -\omega^2 \cos \omega t \cdot V(x)$$

$$\frac{\partial^2}{\partial x^2} u = \cos \omega t \cdot V''(x)$$

one frequency

$$-\rho \omega^2 \cos \omega t \cdot V(x) = P \cos \omega t \cdot V''(x) \quad \text{with } V(0) = V(\pi) = 0$$

- $\frac{\rho \omega^2}{P}$ eigenvalue

$V(x) = \frac{d^2}{dx^2} V(x)$ eigenvalue problem for $\frac{d^2}{dx^2}$

linear differential operator

eigen-form (eigen-functions) $V_k(x) = \sin kx$, $k = 1, 2, 3, \dots$

eigenvalues $k^2 = \frac{\rho \omega^2}{P}$ or eigen-frequencies $\omega_k = \sqrt{\frac{P}{\rho}} \cdot k$

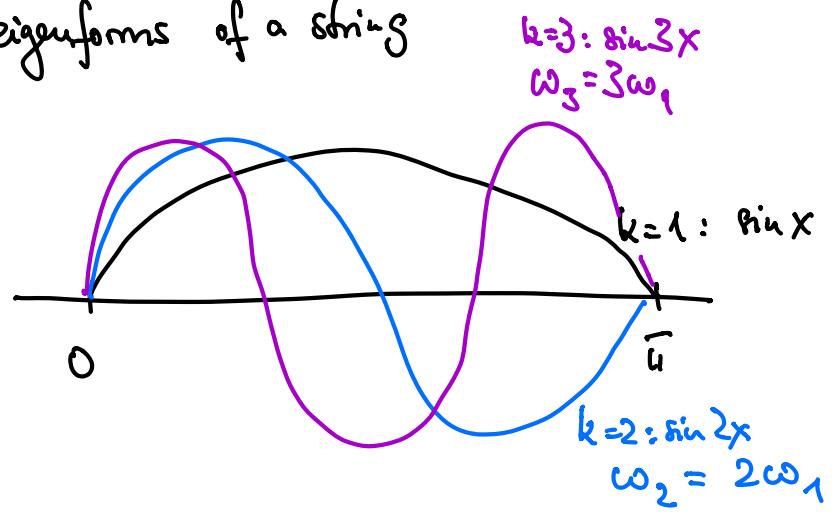
one eigen-oscillation $u_k(t, x) = \cos \omega_k t \cdot \sin kx$

Fourier-series of u is $u(t, x) = \sum_{k=1}^{\infty} g_k \cos \omega_k t \cdot \sin kx$

Superposition / linear combination
of eigen-oscillating

$k=0$

eigenforms of a string



Rem: standard Fourier Series

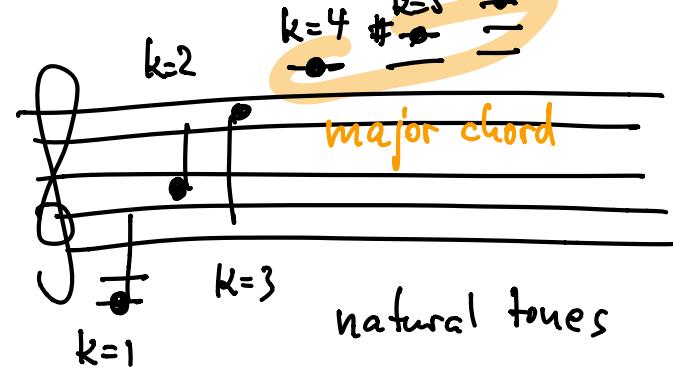
$$f(x) = \frac{a_0}{2} + \sum_{k=1}^{\infty} a_k \cos kx + b_k \sin kx$$

↑ ↑ ↑
 eigen forms eigen forms eigen forms

of an infinitely long string
under tension with a $2\bar{u}$ -periodic
deformation.



Solutions



resp. eigen-value problem

$$\lambda V(x) = \frac{d^2}{dx^2} V(x)$$

for $2\bar{u}$ -periodic and
continuous function V

i.e. boundary condition

$$V(0) = V(2\bar{u}), V'(0) = V'(2\bar{u})$$

$$\begin{aligned} \sin kx, k &= 1, 2, 3, \dots, \lambda = -k^2 \\ \cos kx, k &= 1, 2, 3, \dots, \lambda = -k^2 \\ 1, &\lambda = 0 \end{aligned}$$

Have a good start to your CSE studies.